

Various models around the Cucker-Smale model and their flocking results

Tagung des Deutsch-Französischen Doktorandenkollegs

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Joint work with P. Cattiaux and Laure Pédèches

Modelling collective behavior... Flocking

"Flocking" behavior is a particular kind of collective behavior that can be easily found in nature while observing the collective motion of a large number of individuals.



Modelling collective behavior... Flocking

"Flocking" behavior is a particular kind of collective behavior that can be easily found in nature while observing the collective motion of a large number of individuals.



"Flocking" property means that :

- the distance between two individuals **remains bounded**
- individuals move **in the same direction**

Modelling collective behavior... Flocking

"Flocking" behavior is a particular kind of collective behavior that can be easily found in nature while observing the collective motion of a large number of individuals.



Modelling issues :

- How local interactions at the individuals scale may lead to collective behavior ?
- Which kind of rules drive the local interactions ?

- 1 A few Cucker-Smale models, What about flocking?
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- 2 What about adding noise?
 - What kind of noise? Where?
 - Different notions of stochastic flocking
- 3 Three examples in this framework
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 - Noisy environment and general communication rate
 - Flocking with positive probability

Cucker-Smale Model

Founding papers : Cucker S. et Smale S. 2007

- Cucker F., Smale S., "On the mathematics of emergence", *Japan J. Math.* 2007
- Cucker F., Smale S., "Emergent behavior in flocks", *IEEE Trans. Automat. Control*, 2007

Consider a group of N individuals, the i -th being represented by its position $x_i \in \mathbb{R}^d$ and velocity $v_i \in \mathbb{R}^d$.

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t)(v_j - v_i).$$

- λ measures the strength of the interaction force between individuals.
- Fonction $t \mapsto (\psi_{ij}(t))_{ij}$ is called *communication rate* and $\psi_{ij}(t) \geq 0$ characterises the influence of individual j on individual i .
- A classical choice is $\psi_{ij}(t) = \psi(|x_i(t) - x_j(t)|)$ where ψ is usually chosen positive, decreasing, (ex $\psi(r) = \frac{1}{(1+r^2)^\beta}$, for a given β).

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Cucker-Smale model is thus an **agent-centered** (or microscopic) **mean-field deterministic model, linear on velocities**.

A velocity-attracting model

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi(|x_i - x_j|)(v_j - v_i). \quad (1)$$

Remark : Somming (1) over i , leads to :

- $\bar{v}(t) := \frac{1}{N} \sum_{i=1}^N v_i(t) = \bar{v}^0$ **constant**
- $\bar{x}(t) := \frac{1}{N} \sum_{i=1}^N x_i(t) = \bar{x}^0 + t\bar{v}^0$.

Consider, for $t > 0$

$$z(t) = \sum_{i=1}^N \sum_{j=1}^N |v_i(t) - v_j(t)|^2 \left(= 2N \sum_{i=1}^N |v_i(t) - \bar{v}(t)|^2 \right).$$

We thus have :

$$\frac{dz}{dt} = -\frac{\lambda}{N} \sum_{i=1}^N \sum_{j=1}^N \psi(|x_i - x_j|) |v_i - v_j|^2 \leq 0.$$

Can we prove the alignment of velocities along \bar{v} ? Under which conditions?

Flocking

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t)(v_j - v_i).$$

We say that the group of individuals $\{(x_i(t), v_i(t))\}_{i=1}^N$ *flocks* if :

$$\forall 1 \leq i, j \leq N, \quad \sup_{t \geq 0} |x_i(t) - x_j(t)| < \infty, \quad \lim_{t \rightarrow \infty} |v_i(t) - v_j(t)| = 0$$

The flocking condition can be re-written using the center of mass and the mean velocity :

$$\forall 1 \leq i \leq N, \quad \sup_{t \geq 0} |x_i(t) - \bar{x}(t)| < \infty \quad \lim_{t \rightarrow \infty} |v_i(t) - \bar{v}(t)| = 0$$

or equivalently :

$$\forall 1 \leq i \leq N, \quad \sup_{t \geq 0} |x_i(t) - \bar{x}(t)| < \infty \quad \lim_{t \rightarrow \infty} z(t) = 0$$

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Bounded from below $\psi : \forall r \in \mathbb{R}, |\psi(r)| \geq \ell$

CS'07, bounded from below ψ

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi(|x_i(t) - x_j(t)|)(v_j - v_i).$$

Suppose that ψ is bounded from below by $\ell > 0$, then :

$$\frac{dz}{dt} = -\frac{\lambda}{N} \sum_{i=1}^N \sum_{j=1}^N \psi(|x_i - x_j|) |v_i - v_j|^2 \leq -\frac{\lambda \ell}{N} z(t)$$

thus : $z(t) \leq z^0 e^{-\frac{\lambda \ell}{N} t} \xrightarrow[t \rightarrow \infty]{} 0$ and, for all $t \geq 0$ and all $1 \leq i, j \leq N$:

$$|x_i(t) - x_j(t)| \leq |x_i^0 - x_j^0| + \int_0^t \sqrt{z(s)} ds \leq |x_i^0 - x_j^0| + \frac{2\sqrt{z^0} N}{\lambda \ell} \text{ bounded.}$$

NB : If z decreases fast enough, then, the group necessarily swarms.

Case $\psi(r) = \frac{1}{(1+r^2)^\beta}$, *a priori* not bounded from below

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \frac{1}{(1 + |x_i(t) - x_j(t)|^2)^\beta} (v_j - v_i).$$

Let

$$\psi_\ell(u) = \inf_{0 \leq r \leq u} \psi(r) \text{ and } T_R = \inf \left\{ t, \max_{1 \leq i, j \leq N} |x_i(t) - x_j(t)| \geq R \right\},$$

then, for all $t \leq T_R$ and $1 \leq i, j \leq N$:

$$|x_i(t) - x_j(t)| \leq |x_i^0 - x_j^0| + \frac{\sqrt{z(0)}N}{\lambda\psi_\ell(R^2)}.$$

There thus exists initial data that lead to flocking, whatever β is.

A few flocking results... case $\psi(r) = \frac{1}{(1+r^2)^\beta}$

[Cucker Smale '07], [Ha, Tadmor'08], [Ha Liu '09]

Let $(x_i(t), v_i(t))_{1 \leq i \leq N}$ be the solution to (1) associated with initial data $(x_i^0, v_i^0)_{1 \leq i \leq N}$,

"Unconditional flocking" : case $\beta \in [0, 1/2]$

There exist $x_m > 0$ and $x_M > 0$ such that, for all $t \geq 0$, $1 \leq i \leq N$:

$$x_m \leq |x_i(t) - \bar{x}(t)| \leq x_M, \quad \text{et } |v_i(t) - \bar{v}| \leq |v_i^0 - \bar{v}| e^{-\psi_M t}$$

"Conditional flocking" : case $\beta > 1/2$

If moreover $(x_i^0, v_i^0)_{1 \leq i \leq N}$ satisfy

$$(1 + 2N \|x^0 - \bar{x}^0\|)^{\frac{1-2\beta}{2}} > \frac{3N(2N)^{3/2}}{\lambda} \|v^0 - \bar{v}^0\| \left[\left(\frac{1}{2\beta}\right)^{\frac{1}{2\beta-1}} - \left(\frac{1}{2\beta}\right)^{\frac{1-2\beta}{2}} \right],$$

Then conclusion still holds.

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Non-Symmetric case : Which difference ?

$$\frac{dx_i}{dt} = v_i, \quad \frac{dv_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(X(t))(v_j - v_i)$$

The communication rate is said to be *non symmetric* when

$$\psi_{ij}(X(t)) \neq \psi_{ji}(X(t)).$$

Cucker-Smale '07 : $\psi_{ij}(X(t)) = \frac{1}{(1+|x_i(t)-x_j(t)|^2)^\beta}$ symmetric

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Principal difference : Summing (1) over i used to lead to :

- $\frac{d\bar{v}}{dt} = 0$

- the equation was *dissipative* :

$$\frac{d}{dt} \left(\sum_{i,j} |v_i(t) - v_j(t)|^2 \right) = -\frac{\lambda}{N} \sum_{i,j} \psi_{ij}(X(t)) |v_i(t) - v_j(t)|^2 \leq 0$$

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What kind of noise ? Where ?

Goal : Add noise into the fully deterministic Cucker-Smale interaction.

Questions ?

- Introduce a stochastic term into the kinetic mean-field dynamic (diffusion term ? which form ?)
- Define a stochastic counterpart for the flocking property
- Asymptotic time behavior of the stochastic Cucker-Smale-inspired models ?

What kind of noise? Where?

Personal freedom...

- Each individual has its own alea
- Modelled by a diffusion term of form $\sigma_i(t)dW_i(t)$
 - ▶ with W_i independent d-dimensional Brownian motions
 - ▶ with $\sigma_i(t)$ only depending on $(x_i(t), v_i(t))$
- [Cucker, Mordecki'08], [Ha, Lee, Levy'09] : $\sigma_i = \sqrt{D}I_d$, [Pédèches'16]

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_j(t) - v_i(t)) dt + \sigma(x_i(t), v_i(t)) dW_i(t)$$

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Noisy environnement...

- Common noise for all the individuals, intensity might depend on position/velocity of each individual
- Modelled by a diffusion term of form $\sigma_i(t)dW(t)$, with W d-dimensional Brownian motion
- [Ahn, Ha'10]

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_j(t) - v_i(t)) dt + \sigma(x_i(t), v_i(t)) dW(t)$$

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Noisy perception in the interaction...

- Imperfect perception of the distance with the others
- Modeled by a diffusion term of form $\sum_{j=1}^N \sigma_{i,j}(t)(v_j(t) - v_i(t))dW_{i,j}(t)$, with $W_{i,j}$ d-dimensional independent Brownian motions
- [Ton, Link, Yagi'14], [Erban, Haskovec, Sun '15], [Sun-Lin '15]

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N (v_j(t) - v_i(t)) [\psi_{ij}(t) dt + \sigma_{ij}(t) dW_{ij}(t)]$$

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Which differences ?

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_j(t) - v_i(t)) dt + \sigma(x_i(t), v_i(t)) dW_i(t) \quad (2)$$

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_j(t) - v_i(t)) dt + \sigma(x_i(t), v_i(t)) dW(t) \quad (3)$$

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N (v_j(t) - v_i(t)) [\psi_{ij}(t) dt + \sigma_{ij}(t) dW_{ij}(t)] \quad (4)$$

Equilibrium : Remember that the CS model (1) admits $v_i(t) = \bar{v}^0$ as an equilibrium of the velocities ($\psi_{ij} = \psi_{ji}$)

- In the case (4) it's still true,
- In the general case of models (2) and (3), there is no immediate equilibrium.

Case of a noisy environnement...

In the case (3), if $\sigma_i(t) = D(v_i(t) - v_e)$ where $v_e \in \mathbb{R}^d$ is given, then $v_i = v_e$ is an equilibrium (see [Ahn, Ha '10])

$$dv_i(t) = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_j(t) - v_i(t)) dt + \sigma(x_i(t), v_i(t)) dW(t)$$

Case σ constant : the dynamics can be split into two parts :

- the dynamics of the mean velocity : $\bar{v}(t)$ is driven by a purely stochastic process : $d\bar{v}(t) = \sigma dW(t)$
- the distance to the mean velocity : $\hat{v}_i(t) = v_i(t) - \bar{v}(t)$ satisfies the initial deterministic Cucker-Smale problem :

$$\frac{d\hat{v}_i}{dt} = \frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (\hat{v}_j(t) - \hat{v}_i(t))$$

And now... Laure...

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How to define stochastic flocking? [Cattiaux, D., P. '17]

Recall that **deterministic flocking** is defined as, for all $i \in \{1, \dots, N\}$,

$$\lim_{t \rightarrow \infty} |v_i(t) - \bar{v}(t)| = 0 \quad \text{and} \quad \sup_{0 \leq t < \infty} |x_i(t) - \bar{x}(t)| < \infty ;$$

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The most natural forms of **random flocking** :

- **almost-sure flocking** : the definition above holds **almost surely** :

$$\lim_{t \rightarrow \infty} |v_i(t) - \bar{v}(t)| = 0 \text{ a.s.} \quad \text{and} \quad \sup_{0 \leq t < \infty} |x_i(t) - \bar{x}(t)| < \infty \text{ a.s.};$$

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- **$\mathbb{L}^{p,q}$ -flocking** : **convergence in \mathbb{L}^p of the velocities** towards the center of mass, **boundedness of the positions around their center of mass in \mathbb{L}^q** :

$$\lim_{t \rightarrow \infty} \mathbb{E} [|v_i(t) - \bar{v}(t)|^p] = 0 \quad \text{and} \quad \sup_{0 \leq t < \infty} \mathbb{E} [|x_i(t) - \bar{x}(t)|^q] < \infty.$$

If $q = 1$, we simply say that there is \mathbb{L}^p -flocking.

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If $q = 1$, we simply say that there is \mathbb{L}^p -flocking.

There are others : **mean flocking** and **weak flocking**, but these two are the most demanding.

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Noisy environment and constant communication rate

Take the **same d -dimensional random noise** W impacting all particles :
for $i \in \{1, \dots, N\}$, $k \in \{1, \dots, d\}$,

$$dv_i^k(t) = -\lambda\psi \left(v_i^k(t) - \bar{v}^k(t) \right) dt + D \left(v_i^k(t) - v_e^k \right) dW^k(t)$$

with

- $\psi > 0$ constant communication rate ;
- $W = (W^1, \dots, W^d)$, W^k a Brownian motion ;
- $D > 0$ and $v_e \in \mathbb{R}^d$.

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General strategy :

- Step 1 : study of the evolution of $\bar{v}(t)$;
- Step 2 : study of the distance to the mean : $\hat{v}_i(t) = v_i(t) - \bar{v}(t)$.

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Step 1 : macroscopic scale $d\bar{v}^k(t) = D(\bar{v}^k(t) - v_e^k)dW^k(t)$ and thus :

$$\bar{v}^k(t) = v_e^k + (\bar{v}^k(0) - v_e^k)e^{DW_t^k - \frac{D^2 t}{2}} \xrightarrow[t \rightarrow +\infty]{} v_e^k \text{ p.s.}$$

Noisy environment and constant communication rate

Take the **same d -dimensional random noise** W impacting all particles :
for $i \in \{1, \dots, N\}$, $k \in \{1, \dots, d\}$,

$$dv_i^k(t) = -\lambda\psi \left(v_i^k(t) - \bar{v}^k(t) \right) dt + D \left(v_i^k(t) - v_e^k \right) dW^k(t)$$

Step 2 : microscopic scale $d\hat{v}_i^k(t) = -\lambda\psi\hat{v}_i^k(t)dt - D\hat{v}_i^k dW^k(t)$, hence

$$\hat{v}_i^k(t) = \hat{v}_i^k(0)e^{DW_t^k - (\frac{D^2}{2} + \lambda\psi)t} \xrightarrow[t \rightarrow +\infty]{} 0 \quad p.s.$$

Summary : $\forall i \in \{1, \dots, N\}$, $\forall k \in \{1, \dots, d\}$,

$$v_i^k(t) = \bar{v}^k(t) + \hat{v}_i^k(t) \xrightarrow[t \rightarrow +\infty]{} v_e^k \quad p. s.$$

\Rightarrow **Unconditional almost sure flocking.**

What about the other kinds of stochastic flocking?

Noisy environment and constant communication rate

$$\bar{v}^k(t) = v_e^k + (\bar{v}^k(0) - v_e^k)e^{DW_t^k - \frac{D^2}{2}t}$$

$$\hat{v}_i^k(t) = \hat{v}_i^k(0)e^{DW_t^k - (\frac{D^2}{2} + \lambda\psi)t}$$

\mathbb{L}^1 -flocking :

$$\mathbb{E} \left(|\hat{v}_i^k(t)| \right) = |\hat{v}_i^k(0)| e^{-\lambda\psi t} \mathbb{E} \left(e^{DW_t^k - \frac{D^2}{2}t} \right) = |\hat{v}_i^k(0)| e^{-\lambda\psi t} \xrightarrow{t \rightarrow +\infty} 0.$$

Positions : $\bar{x}_i^k(t) = \bar{x}_i^k(0) + \int_0^t \bar{v}_i^k(s) ds$.

Hence :

$$\begin{aligned} \sup_{t \geq 0} \mathbb{E} \left(|\hat{x}_i^k(t)| \right) &\leq |\hat{x}_i^k(0)| + \int_0^{+\infty} \mathbb{E} \left(|\hat{v}_i^k(s)| \right) ds \\ &\leq |\hat{x}_i^k(0)| + \int_0^{+\infty} |\hat{v}_i^k(0)| e^{-\lambda\psi s} ds < \infty \end{aligned}$$

\Rightarrow Unconditional \mathbb{L}^1 -flocking.

Noisy environment and constant communication rate

$$\bar{v}^k(t) = v_e^k + (\bar{v}^k(0) - v_e^k)e^{DW_t^k - \frac{D^2 t}{2}}$$

$$\hat{v}_i^k(t) = \hat{v}_i^k(0)e^{DW_t^k - (\frac{D^2}{2} + \lambda\psi)t}$$

\mathbb{L}^2 -flocking :

Positions : as before,

$$\sup_{t \geq 0} \mathbb{E} \left(|\hat{x}_i^k(t)| \right) \leq |\hat{x}_i^k(0)| + \int_0^{+\infty} |\hat{v}_i^k(0)| e^{-\lambda\psi s} ds < \infty$$

Velocities :

$$\mathbb{E} \left(|\hat{v}_i^k(t)|^2 \right) = |\hat{v}_i^k(0)|^2 e^{(D^2 - 2\lambda\psi)t} \mathbb{E} \left(e^{2DW_t^k - 2D^2 t} \right) = |\hat{v}_i^k(0)|^2 e^{(D^2 - 2\lambda\psi)t}.$$

\mathbb{L}^2 -flocking $\Leftrightarrow D^2 < 2\lambda\psi$

A case with almost sure flocking, but no \mathbb{L}^2 -flocking

Two realizations of $t \mapsto |\hat{v}(t)| = |v(t) - \bar{v}(t)|$, with an initial configuration for which there is **no \mathbb{L}^2 -flocking**, with

- $d = 2$;
- $N = 9$;
- $\lambda = 10$;
- $D = 7$;

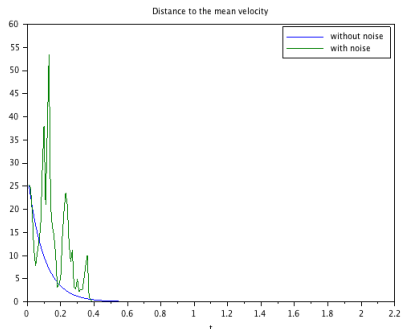
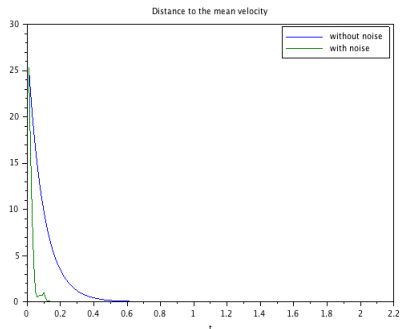


FIGURE: Evolution of $t \mapsto |\hat{v}(t)| = |v(t) - \bar{v}(t)|$.

- 1 A few Cucker-Smale models, What about flocking?
 - Cucker-Smale Model 2007, Flocking
 - Choice of a symmetric communication rate
 - What about non-symmetric comm. rates?
- 2 What about adding noise?
 - What kind of noise? Where?
 - Different notions of stochastic flocking
- 3 Three examples in this framework
 - Noisy environment and constant communication rate
 - Noisy environment and general communication rate
 - Flocking with positive probability

Noisy environment and general communication rate

Take the **same d -dimensional random noise** $W(t)$ impacting all particles : for $i \in \{1, \dots, N\}$, $k \in \{1, \dots, d\}$

$$dv_i^k = -\frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_i^k - v_j^k) dt + D (v_i^k - v_e^k) dW(t),$$

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with

- $\psi_{ij} = \psi_{ij}(v(\cdot), x(\cdot))$ locally Lipschitz, non-negative and symmetric (for instance the **Cucker-Smale rate**);
- $D > 0$ and $v_e \in \mathbb{R}^d$.

Theorem [Cattiaux-D.-P. '17]

The system **flocks almost surely**. However, if $2\lambda \sup_{i,j,x,v} \psi_{ij}(v, x) \leq D^2$, there is **no \mathbb{L}^2 -flocking**.

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Theorem [Cattiaux-D.-P. '17]

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Remarks :

- Struggle between λ and D .
- In the deterministic case, **unconditional flocking** for the Cucker-Smale communication rate **only if** $\gamma \leq 1/2$... **multiplicative noise** "improves" that.

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Flocking with a positive probability

Take the **same d -dimensional random noise** $W(t)$ for all particles,

$$dv_i = -\frac{\lambda}{N} \sum_{j=1}^N \psi_{ij}(t) (v_i - v_j) dt + \sigma(v_i) dW(t),$$

with

- $\psi_{i,j} = \tilde{\psi}(|x_i - x_j|)$ locally Lipschitz, non-negative and non-increasing (for instance the **Cucker-Smale rate**);
- σ globally K -Lipschitz continuous.

Flocking with a positive probability

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- σ globally K -Lipschitz continuous.

Theorem [Cattiaux-D.-P. '17]

Under some assumptions on $x(0)$, $v(0)$, K and $\tilde{\psi}$, there is **flocking with a positive probability**, that is there exists $p \in (0, 1]$ such that

$$\mathbb{P}\left(\forall i \in \{1, \dots, N\}, \lim_{t \rightarrow \infty} |v_i(t) - \bar{v}(t)| = 0\right.$$

$$\left. \text{and } \sup_{0 \leq t < \infty} |x_i(t) - \bar{x}(t)| < \infty\right) \geq p$$

A case with partial flocking

Two realizations of $t \mapsto |\hat{v}(t)| = |v(t) - \bar{v}(t)|$, with an initial configuration – **that does not satisfy the hypotheses of the theorem** – for which there is **flocking in the deterministic case**, with

- $d = 2$;
- $N = 9$;
- $\lambda = 10$;
- σ diagonal,
- $\sigma^{k,k}(v) = 1 + \sin(v^k)$
- $\psi(x, y) = \frac{1}{1 + |x - y|^2}$.

$$dv_i^k(t) = -\frac{10}{9} \sum_{j=1}^9 \frac{v_i^k - v_j^k}{1 + |x_i - x_j|^2} dt + (1 + \sin(v_i^k)) dW^k(t), \quad k \in \{1, 2\}$$

A case with partial flocking

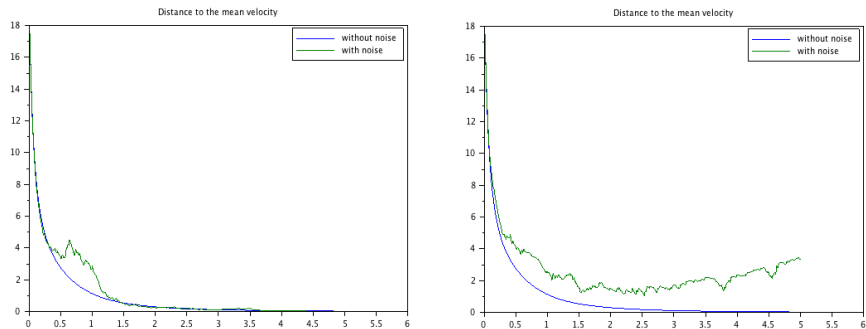


FIGURE: Evolution of $t \mapsto |\hat{v}(t)| = |v(t) - \bar{v}(t)|$.

The **probability of flocking** looks to be **strictly between 0 and 1**...

Thank you !