

# Mass-like invariants for asymptotically hyperbolic manifolds

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# Plan

- 1 Asymptotically hyperbolic manifolds
- 2 The case  $\tau = n$
- 3 Asymptotic invariants for  $\tau > 0$

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# Why asymptotically hyperbolic?

- Cauchy surfaces for asymptotically anti-de Sitter ( $adS^{n+1}$ ) spacetimes,
- Most natural boundary condition “at infinity” after asympt. Euclidean?
- adS/CFT correspondence (conformally compact initial data).

# Definitions

Roughly speaking:  $(M, g, k)$  is *asymptotically hyperbolic* if

$$g \longrightarrow b \text{ and } k \longrightarrow 0 \text{ as } r \rightarrow +\infty,$$

where  $b$  is the metric of  $\mathbb{H}^n$ ,  $b = dr^2 + \sinh^2 r d\Omega_{n-1}^2$ .

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- ... such that  $|\Phi_*g - b|_b(x) = O(e^{-\tau r})$  (and derivatives up to order 2).



# Dependence on the chart?

Proposition (Graham-Lee '93, C-Dahl-Gicquaud '15)

W.l.o.g., one can write

$$\Phi_*g = dr^2 + \sinh^2 r \left[ \sigma_{\mathbb{S}^{n-1}} + \mathbf{m}e^{-\tau r} + O(e^{-(\tau+1)r}) \right] \text{ with}$$
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If  $\Phi_1$  and  $\Phi_2$  are two charts of order  $\tau$ , there exists  $A \in O^+(n, 1)$  s.t.

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- In the Euclidean case:  $\Phi_2 \circ \Phi_1^{-1} = A + O(r^{-(\tau-1)})$ .

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# The mass vector for $\tau = n$ (X. Wang)

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- The *mass vector* is  $\mathbf{p}_\Phi = (p_0, p_1, \dots, p_n) \in \mathbb{R}^{n+1}$  with

$$p_0 = \int_{\mathbb{S}^{n-1}} \text{tr}^\sigma \mathbf{m} \, d\sigma, \quad p_i = \int_{\mathbb{S}^{n-1}} x^i \text{tr}^\sigma \mathbf{m} \, d\sigma, \quad i = 1 \dots n.$$

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If  $\Phi_1$  and  $\Phi_2$  satisfy  $\Phi_2 \circ \Phi_1^{-1} = A + O(e^{-(n+1)r})$  with  $A \in O^+(n, 1)$ , then

$$\mathbf{p}_{\Phi_1} = A \cdot \mathbf{p}_{\Phi_2}.$$

# The mass vector for $\tau = n$ (X. Wang)

Consequences:

- The number  $|\mathbf{p}|_\eta^2 = -p_0^2 + \sum_i p_i^2$  does not depend on  $\Phi$ !
- But  $p_0 = \int_{\mathbb{S}^{n-1}} \text{tr}^\sigma \mathbf{m} \, d\sigma$  alone does depend on  $\Phi$ .
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Example: Kottler metrics for  $n = 3$

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$$g_M := \frac{d\rho^2}{1 - \frac{2M}{\rho} + \rho^2} + \rho^2 \sigma_{\mathbb{S}^2},$$

- We find that  $\mathbf{m} = 3M\sigma$ , thus  $\mathbf{p} = C(M, 0, 0, 0)$ .

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# Asymptotic invariants

## Definition

An asymptotic invariant is a  $O^+(n, 1)$ -intertwining map

$$\Psi : \text{AH}(\tau) \longrightarrow \mathbb{V},$$

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- Example: ( $\tau = n$ ), we have  $\mathbb{V} = \mathbb{R}^{n,1}$ ,  $\Psi(\mathbf{m}) = \mathbf{p}$ .

Results for  $n = 3$ 

## Theorem (C-Dahl-Gicquaud '15)

For all  $\tau \geq 2$ ,  $\tau \in \mathbb{N}$ , there is a unique map  $\Psi_\tau : \Gamma(S^2 T^* \mathbb{S}^{n-1}) \rightarrow \mathbb{V}_\tau$  which is intertwining:

$$\Psi_\tau(\mathbf{m}) = \left( P \in \mathbb{V}_\tau^* \mapsto \int_{\mathbb{S}^2} P(1, x^1, x^2, x^3) \operatorname{tr}^\sigma \mathbf{m} \, d\sigma \right),$$

where  $\mathbb{V}_\tau^* = \{P \in \mathbb{R}[x^0, x^1, x^2, x^3], \deg P = \tau - 2, \square_\eta P = 0\}$ .

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- For  $\tau = 2$ ,  $\mathbb{V}_2 = \mathbb{R}$ , and  $\int_{S^2} \operatorname{tr}^\sigma \mathbf{m} \, d\sigma$  is an asymptotic invariant!

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$$\bar{\Psi}_\tau : g \mapsto \mathbb{V}_\tau^*$$

defined on the set of metrics  $g \in \text{AH}(\frac{\tau}{2} + \epsilon)$  such that  $e^{(\tau-2)r} F_\tau(g)$  is integrable on  $\mathbb{H}^3 \setminus \bar{B}_R$ , where  $F_\tau$  is a geometric differential operator. Moreover, we have  $\bar{\Psi}_\tau(g) = \Psi_\tau(g)$  for  $g \in \text{AH}(\tau)$ .



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- The operators  $F_\tau$  are defined by  $F_\tau(g) := \Delta_g \text{Scal}_g - C(\tau) \text{Scal}_g$  for  $\tau \neq 3$ , and  $F_3 = \text{Scal}$ .

## Open questions

- Case  $n \geq 3$  (Work in progress)
- Complex intertwining maps (in progress)
- Multilinear in  $\mathbf{m}$  intertwining maps (Gauss-Bonnet-Chern masses)

# Representation theory for $O(n, 1)$

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Figure : Nick Hobgood, [http://commons.wikimedia.org/wiki/File:Amphiprion\\_ocellaris\\_%28Clown\\_anemonefish%29\\_Nemo.jpg](http://commons.wikimedia.org/wiki/File:Amphiprion_ocellaris_%28Clown_anemonefish%29_Nemo.jpg)